

5th TLEP Workshop

July 25-26, Fermilab

Top precision measurements at LC - Theory improvements

Markus Schulze



Top quark physics at TLEP

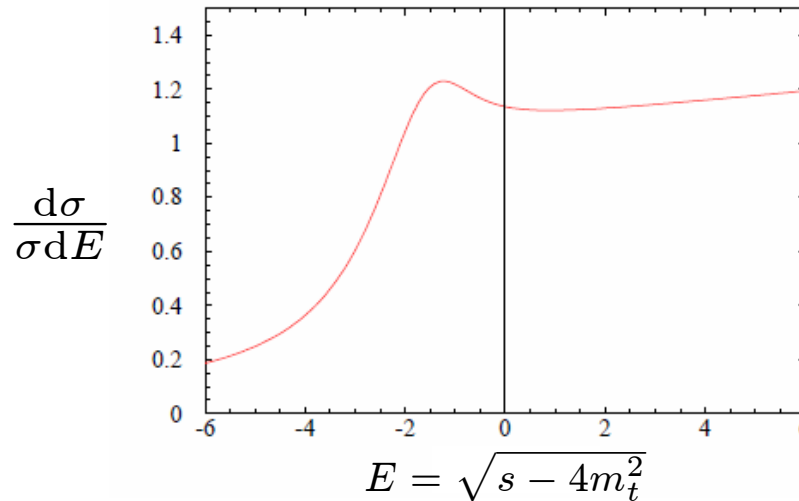
- TLEP at $E=350$ GeV gives us the opportunity to study $t\bar{t}$ production around its threshold.
- The high luminosity and the clean experimental environment at TLEP promise a high precision top quark physics program.
- High precision might be a next logical step if no striking discoveries show up at the LHC, or if a discovery requires further scrutiny.
- Of course, this has to be seen in perspective with prospects for ILC/CLIC and the LHC. First studies indicate that TLEP performance is competitive.
+ It offers a long term vision for HEP in Europe.

Outline

- **Threshold scan**
 - ttbar cross section
 - Top quark mass, width
- **Top quark couplings**
 - Electroweak couplings
 - Yukawa coupling
 - Rare decays, FCNC, single top
- **Light stops**

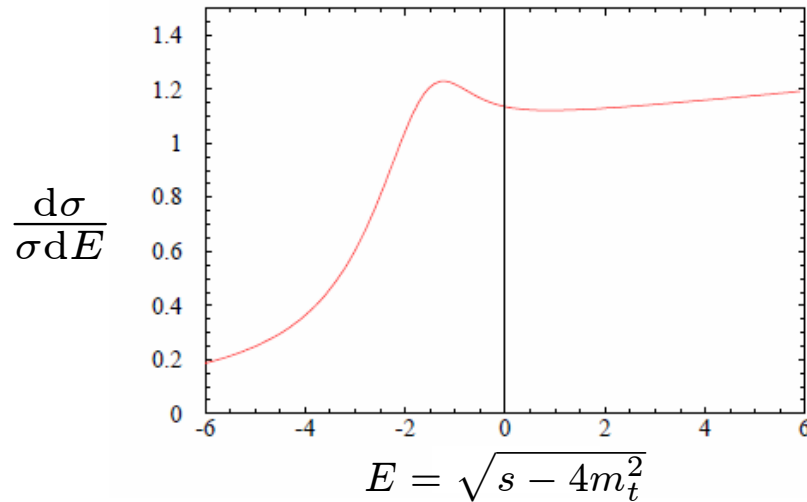
Threshold scan

Threshold scan



- The Rydberg binding energy of toponium $\sim m_t \alpha_s^2 \approx 2$ GeV is similar to the top quark width $\Gamma_t \approx 1.5$ GeV, both of which are much larger than Λ_{QCD} .
- Two opposite effects govern the top quark threshold region:
 - The QCD interaction between the non-relativistic top quarks pulls towards Coulomb-like toponium bound states.
 - The large top quark width $\Gamma_t \gg \Lambda_{\text{QCD}}$ leads to a rapid decay before a bound state can be formed.

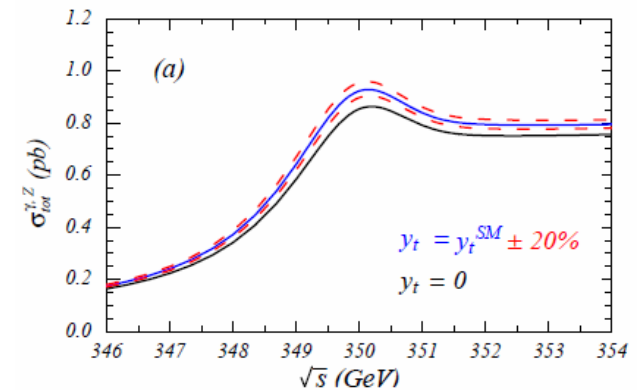
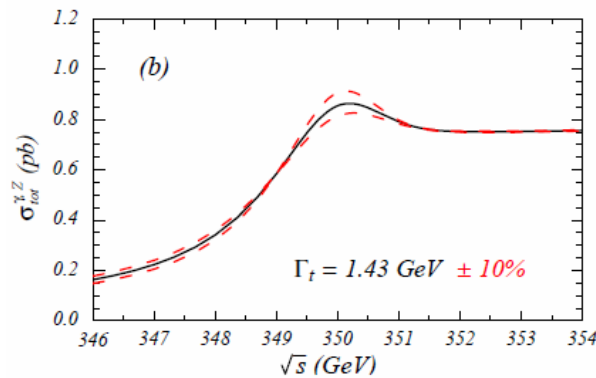
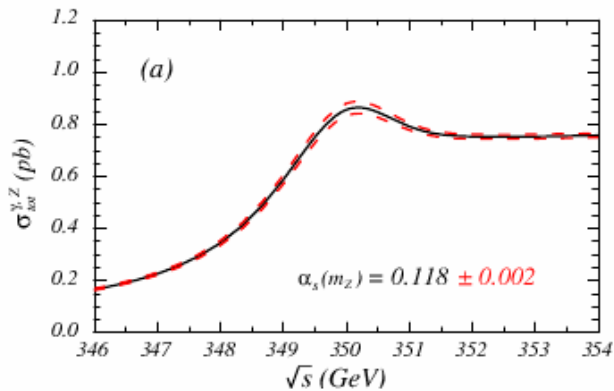
Threshold scan



- The resonance peak is smeared out.
 - The large decay width effectively serves as a cut-off for non-perturbative effects \rightarrow cross section can be described within perturbation theory.
 - The slow velocity close to threshold requires a pert. expansion in α_s^n / v^m
 \rightarrow NRQCD with QCD Coulomb potential.
- \rightarrow This allows a reliable description of the threshold region which is based entirely on first principles.

Threshold scan

[Stewart]

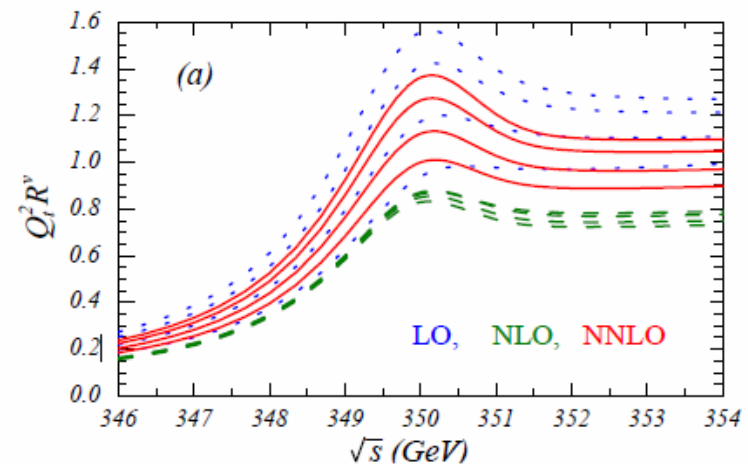
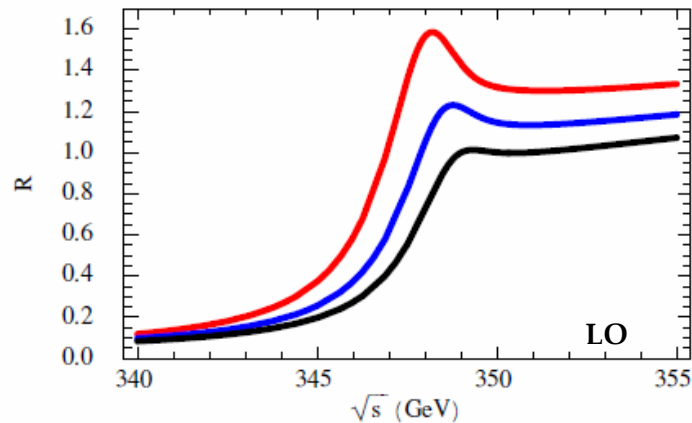


- The resonance cross section $\sigma_{\text{res}} \sim \alpha_s^3 / (m_t \Gamma_t)$ is very sensitive to strong coupling, top quark mass and width.
 - Higgs boson exchange introduces dependence on y_t through loops.
- To what precision can we predict threshold dynamics?
What is the expected experimental sensitivity?

Threshold scan: the ttb cross section

$$R = \frac{\sigma_{tt\bar{b}}}{\sigma_{\mu^+\mu^-}} = v \sum_k \left(\frac{\alpha_s}{v} \right)^k \times \left[1 + \underbrace{\{\alpha_s, v\}}_{\text{LO}} + \underbrace{\{\alpha_s^2, \alpha_s v, v^2\}}_{\text{NLO}} + \dots \right]$$

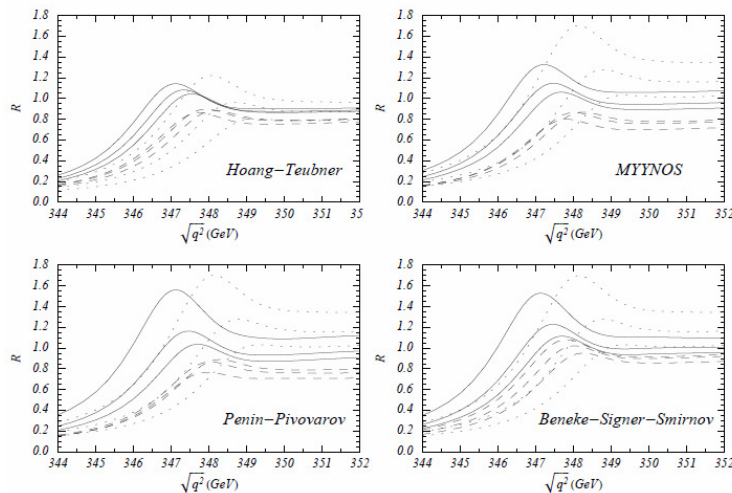
Coulomb resummation



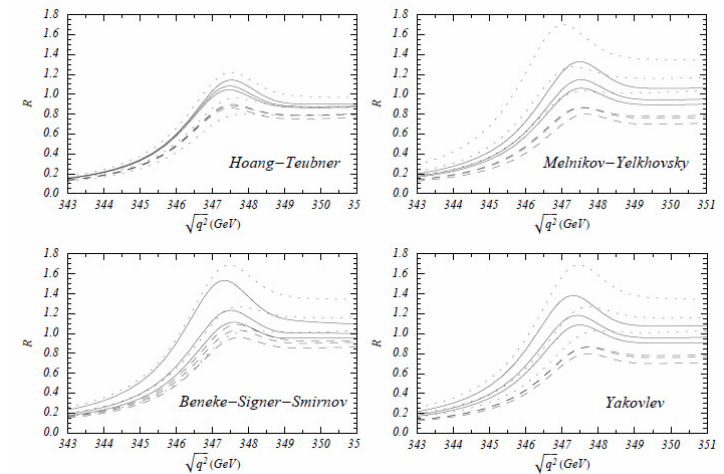
- Description pushed to NNLO by several groups:
[Hoang, Beneke, Melnikov, Nagano, Ota, Penin, Pivovarov, Signer, Smirnov, Sumino, Teubner, Yakovlev, Yelkhovsky]
- Corrections are large, scale variation bands do not overlap
residual scale uncertainty $\sim 20\%$
→ partly contributed to Renormalon contribution of
top quark pole mass, logs of largely different scales ($E_t - v_t - m_t$)

Threshold scan: the ttb cross section

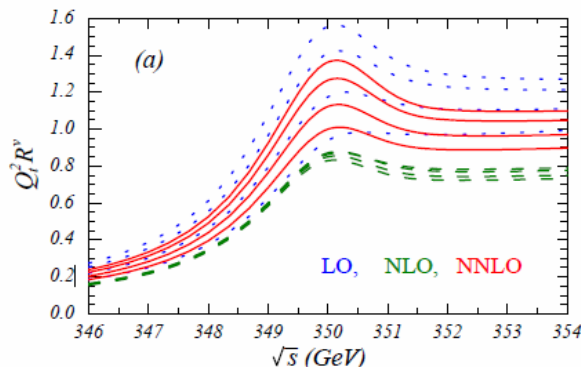
- More appropriate „short-distance“ definitions of top quark mass such as MSbar or threshold mass improve convergence:



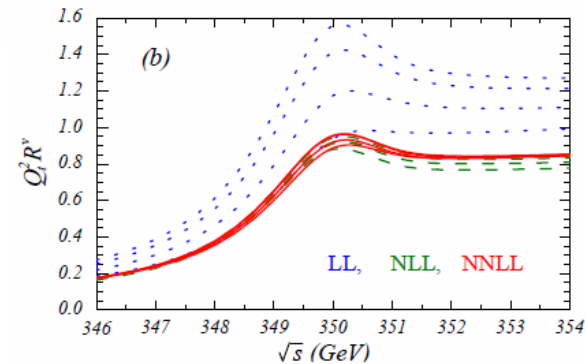
short dist.
masses



- Logarithmic resummation: $R = v \sum_k \left(\frac{\alpha_s}{v} \right)^k \sum_j (\alpha_s \ln v)^j \times \left[1 + \{ \alpha_s, v \} + \{ \alpha_s^2, \alpha_s v, v^2 \} + \dots \right]$



resum
logs(v)

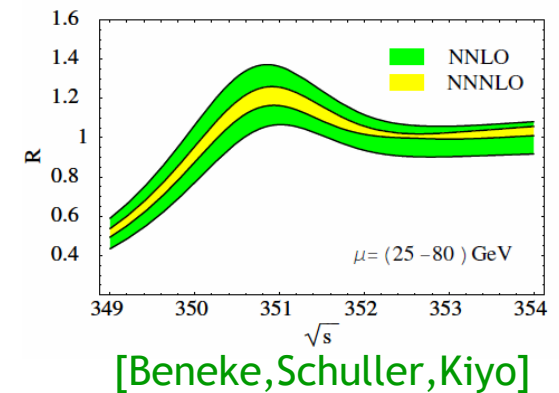


[Stewart]

Threshold scan: the ttb cross section

- Ultimately, for a rigorous quantitative analysis of threshold production the complete N^3LO corrections are required.
- Significant progress towards this goal in recent years. We can expect the full result in the near future.

[Anzei, Beneke, Hoang, Kiyo, Kniehl, Marquard, Penin, Piclum, Schuller, Seidel, Smirnov, Steinhauser, Sumino]

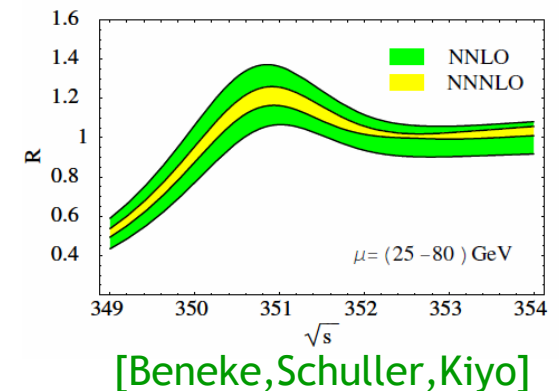


→ The N^3LO uncertainty on total cross section is likely of order 3%.

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- At this level of precision one needs to account for
 - electroweak corrections [Grzadkowski, Kuhn, Krawczyk, Stuart], [Hoang, Reisser]
 - mixed QCD-el.weak corrections [Eiras, Steinhauser], [Kiyo, Seidel, Steinhauser]
 - finite width effects, non-factorizable corrections [Hoang, Reisser, Femenia]

Threshold scan: the top quark mass

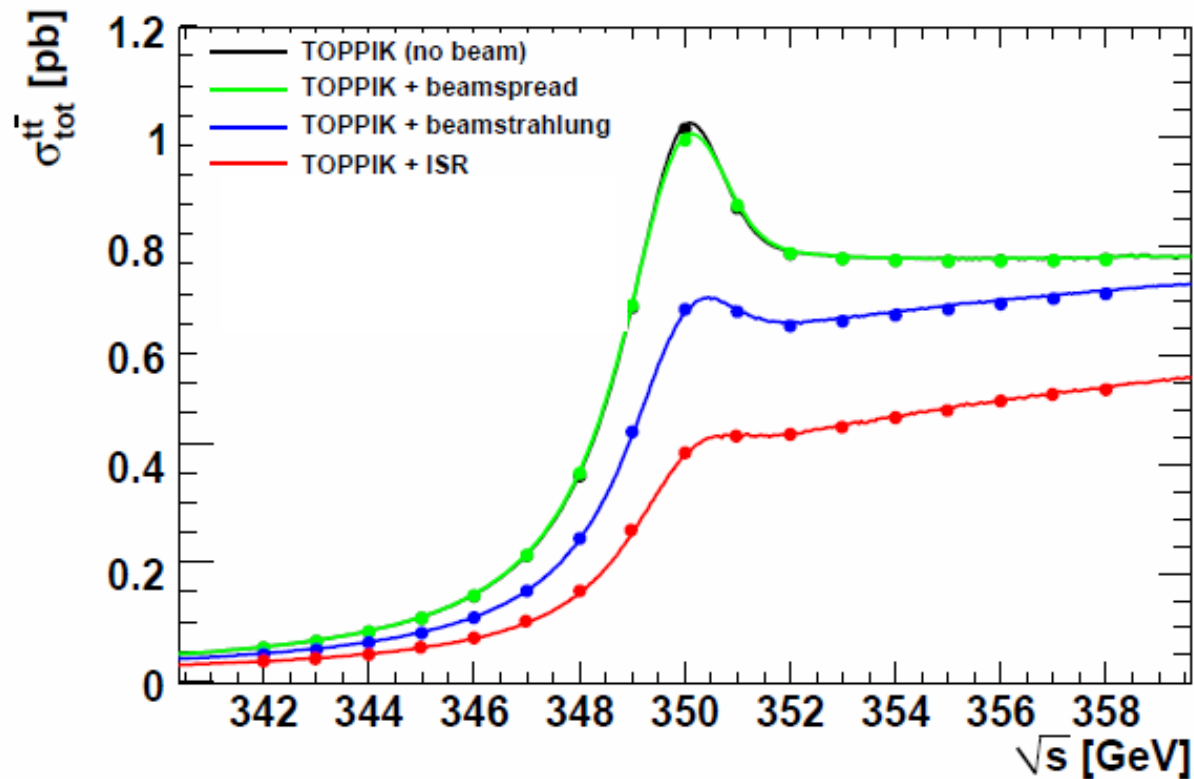
N³QCD relation for resonance energy and m_t : [Penin,Steinhauser]
[Kiyo,Sumino]

$$\sqrt{s_{\text{res}}} = \left[1.9833 + 0.007 \frac{m_t - 174.3 \text{ GeV}}{174.3 \text{ GeV}} \pm 0.0009 \right] \times m_t,$$

- Theoretical uncertainty of **80 MeV** on m_t (pole mass)
- The use of the short-distance MSbar mass reduces uncertainty to **40 MeV**.

Threshold scan: the top quark mass

Realistic studies need to include beam and detector effects, ISR, backgrounds,...



[Gournaris]

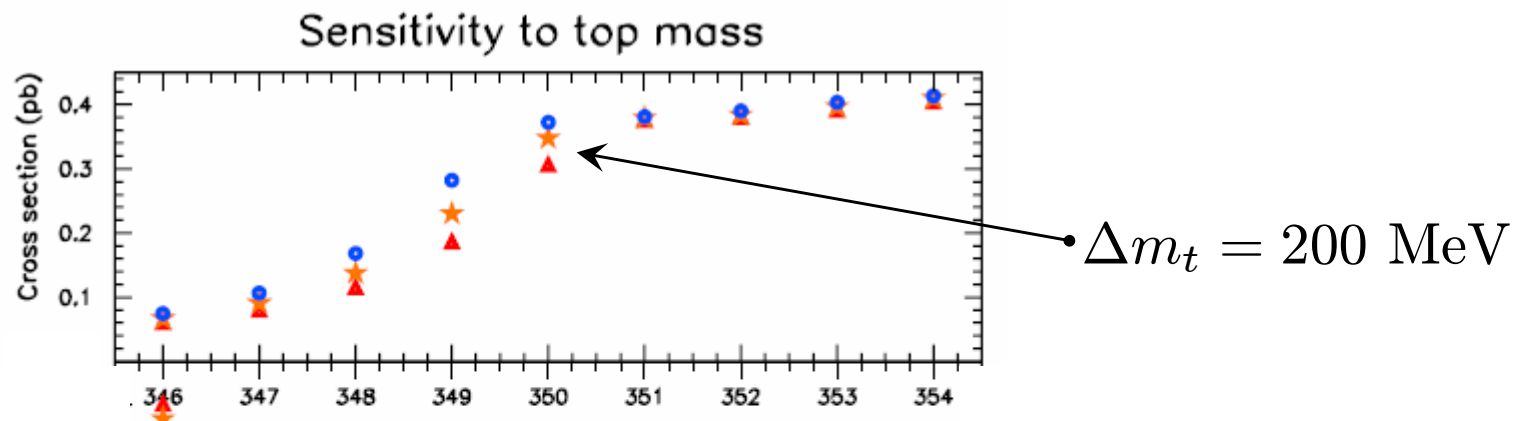
Threshold scan: the top quark mass

[Martinez,Miquel]:

Multi-parameter fits to the $t\bar{t}$ threshold observables at a future e^+e^- linear collider

(Standard reference for a realistic study)

Use NNLO simulation and include detector effects, selection efficiency and backgrounds, assuming the TESLA beam and detector design.



- The simulations show an estimated experimental error of about **3%** on the total cross section (much below the one of the differential observables).
- The resulting uncertainty on m_t is **31 MeV** (from multiparam. fit of 3 observables).
Neglecting uncertainties from beam energy and luminosity spectrum.

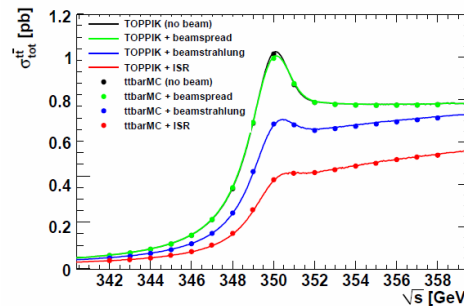
Threshold scan: the top quark mass

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Multi-parameter fits to the $t\bar{t}$ threshold observables at a future e^+e^- linear collider

(Standard reference for a realistic study)

[Gournaris]: Similar study for ILC including beam and luminosity uncertainties



→ In summary, for a 300 fb^{-1} threshold scan the total expected uncertainty on m_t is **100 MeV**, resulting from the sum in quadrature of the following contributions:

31 MeV ([Martinez,Miquel]),

35 MeV (beam energy),

50 MeV (luminosity spectrum) and

80 MeV (from the conversion of s_{res} into m_t).

Above threshold: the top quark mass

$\delta m_t = 100 \text{ MeV}$ in perspective:

- LHC projections: [Snowmass]

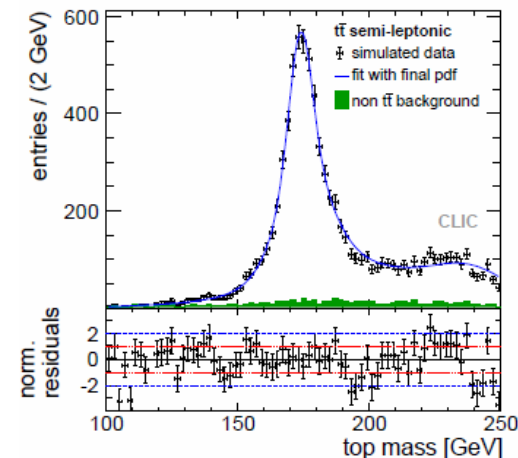
	Projections					Projections			Projections				
CM Energy	14 TeV					14 TeV			14 TeV			33 TeV	100 TeV
Luminosity	$100 fb^{-1}$	$300 fb^{-1}$	$3000 fb^{-1}$	$3000 fb^{-1}$	$3000 fb^{-1}$	$100 fb^{-1}$	$300 fb^{-1}$	$3000 fb^{-1}$	$100 fb^{-1}$	$300 fb^{-1}$	$3000 fb^{-1}$	$3000 fb^{-1}$	$3000 fb^{-1}$
Syst. (GeV)	0.7	0.7	0.6	0.6	0.6	1.0	0.7	0.5	1.5	1.5	1.0	1.0	0.6
Stat. (GeV)	0.04	0.04	0.03	0.03	0.01	0.10	0.05	0.02	1.8	1.0	0.3	0.1	0.1
Total, GeV	0.7	0.7	0.6	0.6	0.6	1.0	0.7	0.5	2.3	1.8	1.1	1.0	0.6

likelihood methods
kinem. end-points
J/Psi method

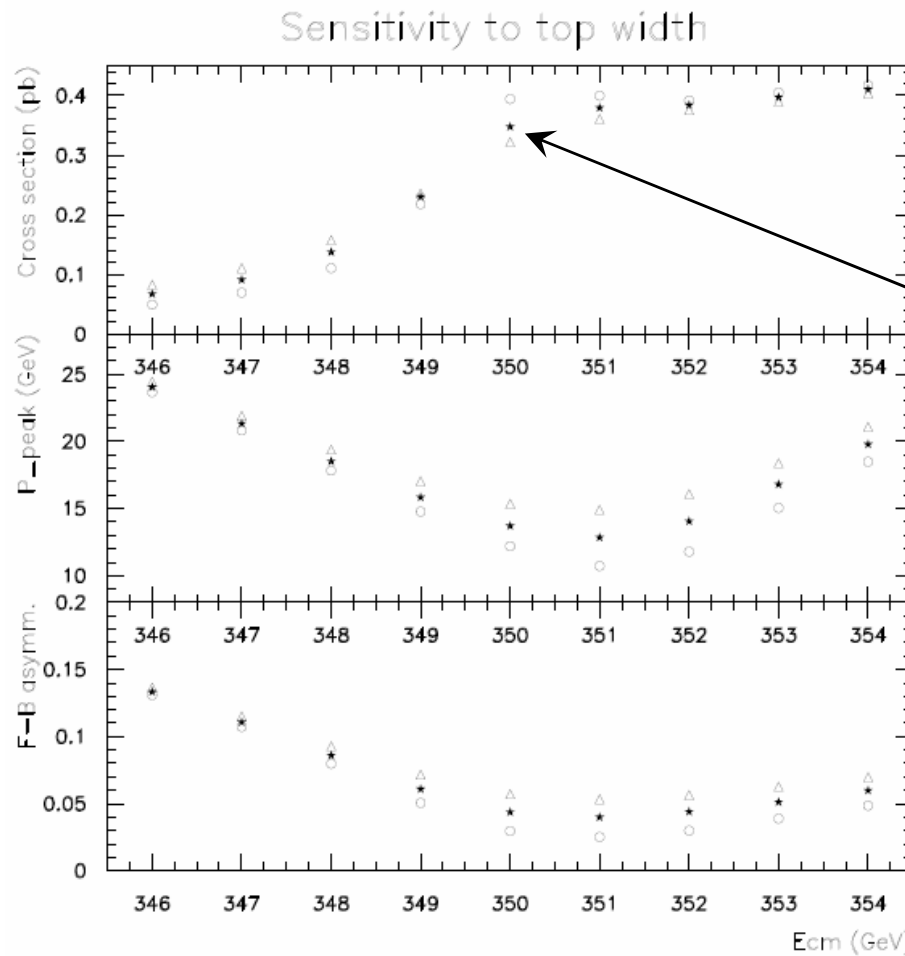
- ILC/CLIC projections: [Seidel, Simon, Tesar]

channel	m_{top}	Δm_{top}	σ_{top}	$\Delta \sigma_{top}$
fully-hadronic	174.049	0.099	1.47	0.27
semi-leptonic	174.293	0.137	1.70	0.40
combined	174.133	0.080	1.55	0.22

Table 2 Results summary for the top mass measurement at 500 GeV for an integrated luminosity of $100 fb^{-1}$. All numbers are given in units of GeV, Errors are statistical only.



Threshold scan: top quark width

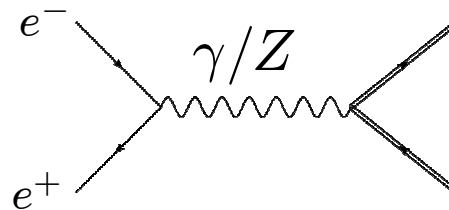
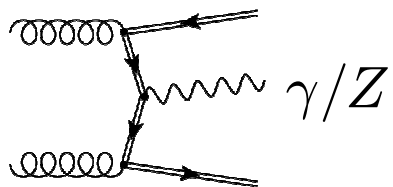


[Martinez,Miquel]

- Accuracy of 2% on top quark width is achievable

Top quark couplings

Top quark electroweak couplings



$$\Gamma_{\mu}^{t\bar{t}V} \sim \gamma_{\mu} (F_{1,V} + \gamma_5 F_{1,A}) + \sigma_{\mu\nu} \frac{q^{\nu}}{2m_t} (iF_{2,V} + \gamma_5 F_{2,A})$$

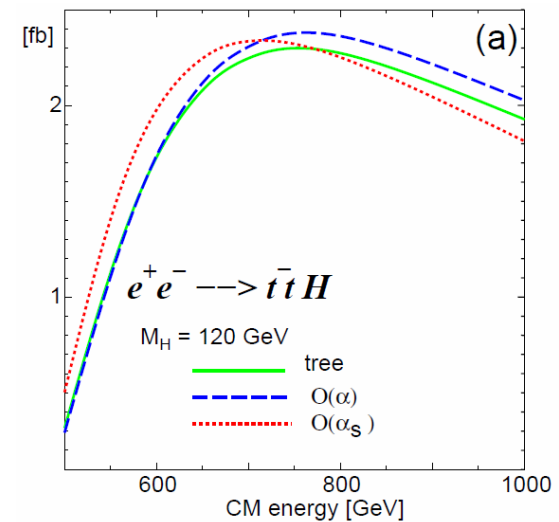
- Currently, these couplings are not very well measured (apart from Q_t).
At hadron colliders only accessible through $t\bar{t}b+\gamma$ and $t\bar{t}b+Z$.
At Tevatron, a handful of $t\bar{t}b+\gamma$ events. Indirect constraints from CLEO, LEP on $t\bar{t}bZ$.
→ Ultimate precision at the LHC (3000 fb^{-1}): **10-20%** [Baur, Juste, Orr, Rainwater]
- In e^+e^- collisions, the cross section has an entangled dependence on photon and Z coupling. Polarized beams and FB asymmetries can be used to get a handle on individual couplings. [ILC TDR], [Devetak, Nomerotski, Peskin]
→ ILC precision with 500 fb^{-1} : **sub percent level** (80:30% polarization) [Baur et al.], [Snowmass reports]
- At TLEP, only vector part accessible (axial couplings require p-wave)
→ no sensitivity to electric dipole moments
most likely compatible with ILC precision (higher lumi but no polarization)

Top quark Yukawa coupling

$$e^+e^- \rightarrow t\bar{t} + H$$

- $E_{\text{thresh}} = 2m_t + m_H = 470 \text{ GeV}$

→ large gain from increased CM energy
 $\sigma(800 \text{ GeV})/\sigma(500 \text{ GeV}) \sim 7$



$\sqrt{s} = 500 \text{ GeV}, \mathcal{L} = 1000\text{fb}^{-1}: \delta y_t/y_t \approx 10\%$

[Dawson,Juste,Reina,Wackerroth]

$\sqrt{s} = 800 \text{ GeV}, \mathcal{L} = 1000\text{fb}^{-1}: \delta y_t/y_t \approx 6 \%$

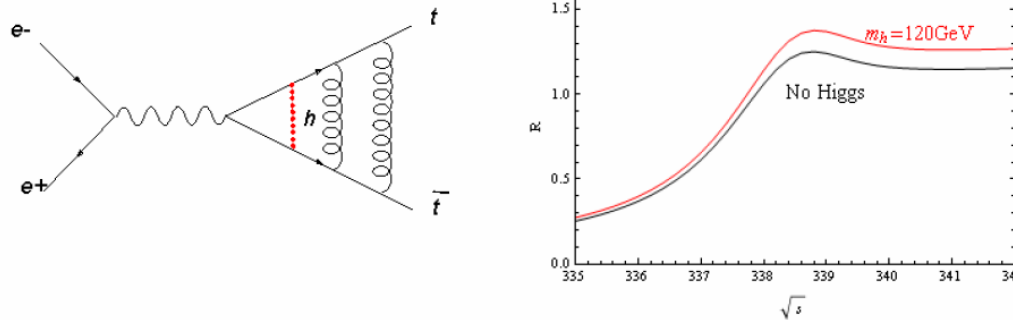
[Yonamine,Ikematsu,Tanabe,
Fujii,Kiyo,Sumino,Yokoya]

- At the LHC, difficult final state with huge background.
Still, 15-20 % precision seems possible.

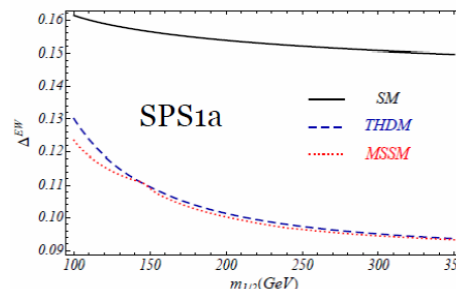
[CMS-NOTE-2012-006],
[ATLAS-PHYS-PUB-2012-004]

Threshold scan: Top quark Yukawa coupling

- At TLEP energies, sensitivity only through loop effects

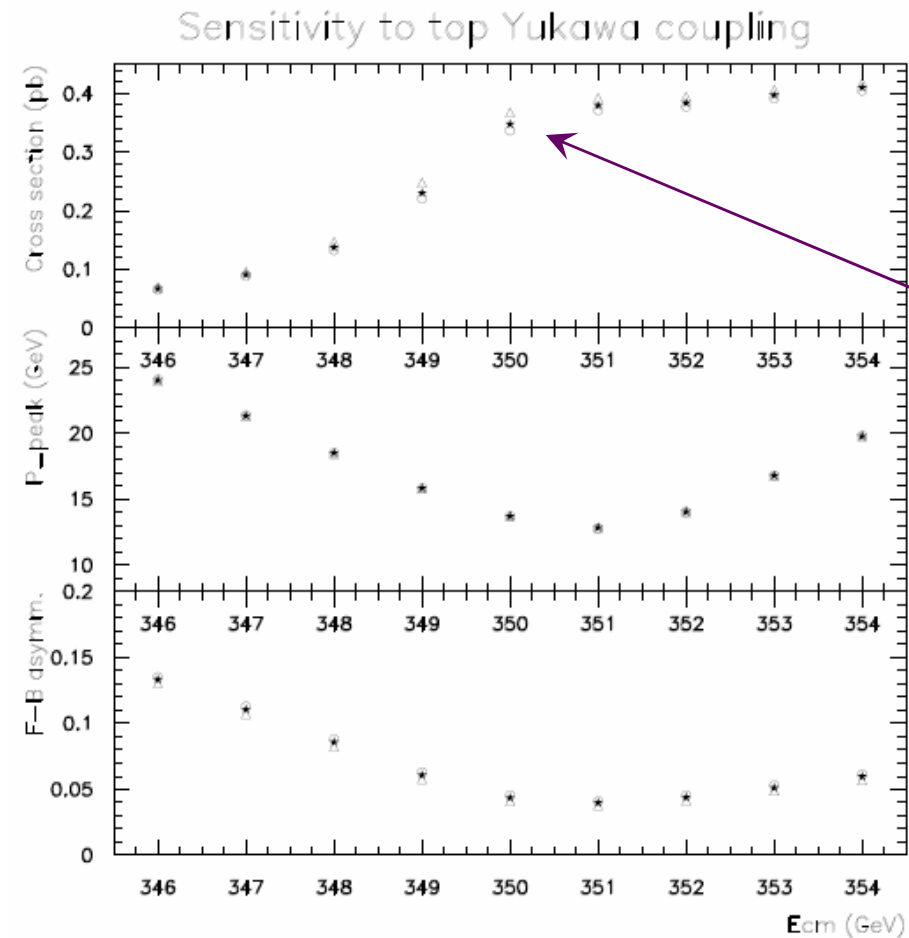


- LO (1-loop) shift: **+6%** (for $M_H = 120 \text{ GeV}$) [Grzadkowski, Kuhn, Krawczyk, Stuart]
- NLO el.weak+QCD (2-loop): **+9%** [Eiras, Steinhauser]
- SM+SUSY NLO corrections: **$\sim 1\%$** (almost complete screening of SM effect)



[Kiyo, Steinhauser, Zerf]

Top quark Yukawa coupling



[Martinez,Miquel]

variation $\delta y_t = \pm 50\%$

- A measurement is challenging but not impossible
- Under optimistic assumptions: 30% accuracy is possible

Rare top quark decays

Rare top quark decays: FCNC

$$e^+e^- \rightarrow t\bar{t} \xrightarrow{\text{FCNC}} (Z/\gamma/g/h + j)(Wb)$$

Process	SM [67]	2HDM(FV) [67, 68]	2HDM(FC) [69]	MSSM [70]	RPV [71, 72]	RS [73, 74]
$t \rightarrow Zu$	7×10^{-17}	–	–	$\leq 10^{-7}$	$\leq 10^{-6}$	–
$t \rightarrow Zc$	1×10^{-14}	$\leq 10^{-6}$	$\leq 10^{-10}$	$\leq 10^{-7}$	$\leq 10^{-6}$	$\leq 10^{-5}$
$t \rightarrow gu$	4×10^{-14}	–	–	$\leq 10^{-7}$	$\leq 10^{-6}$	–
$t \rightarrow gc$	5×10^{-12}	$\leq 10^{-4}$	$\leq 10^{-8}$	$\leq 10^{-7}$	$\leq 10^{-6}$	$\leq 10^{-10}$
$t \rightarrow \gamma u$	4×10^{-16}	–	–	$\leq 10^{-8}$	$\leq 10^{-9}$	–
$t \rightarrow \gamma c$	5×10^{-14}	$\leq 10^{-7}$	$\leq 10^{-9}$	$\leq 10^{-8}$	$\leq 10^{-9}$	$\leq 10^{-9}$
$t \rightarrow hu$	2×10^{-17}	6×10^{-6}	–	$\leq 10^{-5}$	$\leq 10^{-9}$	–
$t \rightarrow hc$	3×10^{-15}	2×10^{-3}	$\leq 10^{-5}$	$\leq 10^{-5}$	$\leq 10^{-9}$	$\leq 10^{-4}$

[Snowmass white paper]

- FCNC top quark decays are highly suppressed in the SM (light quark masses and small CKM angle)
 - New Physics models introduce significantly higher rates
- Any measured deviation from zero indicates NP in the top quark decay

Rare decays: FCNC

[Snowmass white paper]

LHC
projections:

Process	Br Limit	Search	Dataset
$t \rightarrow Zq$	2.2×10^{-4}	ATLAS $t\bar{t} \rightarrow Wb + Zq \rightarrow \ell\nu b + \ell\ell q$	300 fb ⁻¹ , 14 TeV
$t \rightarrow Zq$	7×10^{-5}	ATLAS $t\bar{t} \rightarrow Wb + Zq \rightarrow \ell\nu b + \ell\ell q$	3000 fb ⁻¹ , 14 TeV
$t \rightarrow \gamma q$	8×10^{-5}	ATLAS $t\bar{t} \rightarrow Wb + \gamma q$	300 fb ⁻¹ , 14 TeV
$t \rightarrow \gamma q$	2.5×10^{-5}	ATLAS $t\bar{t} \rightarrow Wb + \gamma q$	3000 fb ⁻¹ , 14 TeV
$t \rightarrow gu$	4×10^{-6}	ATLAS $qg \rightarrow t \rightarrow Wb$	300 fb ⁻¹ , 14 TeV
$t \rightarrow gu$	1×10^{-6}	ATLAS $qg \rightarrow t \rightarrow Wb$	3000 fb ⁻¹ , 14 TeV
$t \rightarrow gc$	1×10^{-5}	ATLAS $qg \rightarrow t \rightarrow Wb$	300 fb ⁻¹ , 14 TeV
$t \rightarrow gc$	4×10^{-6}	ATLAS $qg \rightarrow t \rightarrow Wb$	3000 fb ⁻¹ , 14 TeV
$t \rightarrow hq$	2×10^{-3}	LHC $t\bar{t} \rightarrow Wb + hq \rightarrow \ell\nu b + \ell\ell qX$	300 fb ⁻¹ , 14 TeV
$t \rightarrow hq$	5×10^{-4}	LHC $t\bar{t} \rightarrow Wb + hq \rightarrow \ell\nu b + \ell\ell qX$	3000 fb ⁻¹ , 14 TeV
$t \rightarrow hq$	5×10^{-4}	LHC $t\bar{t} \rightarrow Wb + hq \rightarrow \ell\nu b + \gamma\gamma q$	300 fb ⁻¹ , 14 TeV
$t \rightarrow hq$	2×10^{-4}	LHC $t\bar{t} \rightarrow Wb + hq \rightarrow \ell\nu b + \gamma\gamma q$	3000 fb ⁻¹ , 14 TeV

ILC
projections:

Process	Br Limit	Search	Dataset
$t \rightarrow Zq$	$1.6 (1.7) \times 10^{-3}$	ILC $t\bar{t}, \gamma_\mu (\sigma_{\mu\nu})$	500 fb ⁻¹ , 500 GeV
$t \rightarrow \gamma q$	1.0×10^{-4}	ILC $t\bar{t}$	500 fb ⁻¹ , 500 GeV

- Limits from 500 fb⁻¹ ILC and 3000 fb⁻¹ LHC are compatible
- Expect TLEP to perform even better, given the higher luminosity

Rare decays: single top

$$e^+e^- \xrightarrow[Z/\gamma \text{ coupl.}]{\text{anomal.}} tq$$

- Studies are possible at $E=250$ GeV (maximal cross section) and at $E=500$ GeV (lower background)

[Aguilar-Saavedra,Riemann]
[Snowmass white paper]

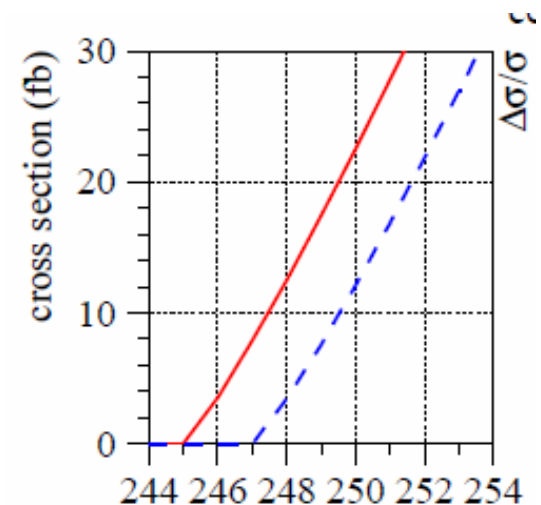
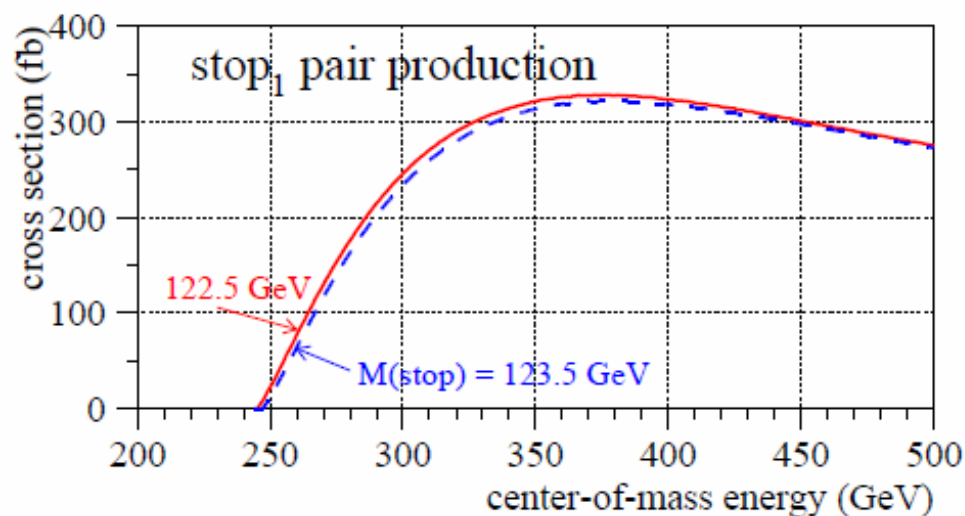
Process	Br Limit	Search	Dataset
$t \rightarrow Zq$	$5(2) \times 10^{-4}$	ILC single top, $\gamma_\mu (\sigma_{\mu\nu})$	500 fb ⁻¹ , 250 GeV
$t \rightarrow Zq$	$1.5(1.1) \times 10^{-4} (-5)$	ILC single top, $\gamma_\mu (\sigma_{\mu\nu})$	500 fb ⁻¹ , 500 GeV
$t \rightarrow \gamma q$	6×10^{-5}	ILC single top	500 fb ⁻¹ , 250 GeV
$t \rightarrow \gamma q$	6.4×10^{-6}	ILC single top	500 fb ⁻¹ , 500 GeV

- Possibly stronger limits than from top quark pair production
- Opportunities below ttb threshold ($E=240..250$ GeV) are not well studied (above results rely on extrapolations)

Light stops

New Physics at around 350 GeV: light stops

- Heavy stop quark scenarios are mostly covered by LHC searches
- Light stops and compressed spectra are difficult, even for HL-LHC
- At lepton colliders: $e^+e^- \rightarrow \tilde{t}_1 \tilde{t}_1^* \rightarrow c\tilde{\chi}_1^0 \bar{c}\tilde{\chi}_1^0$.



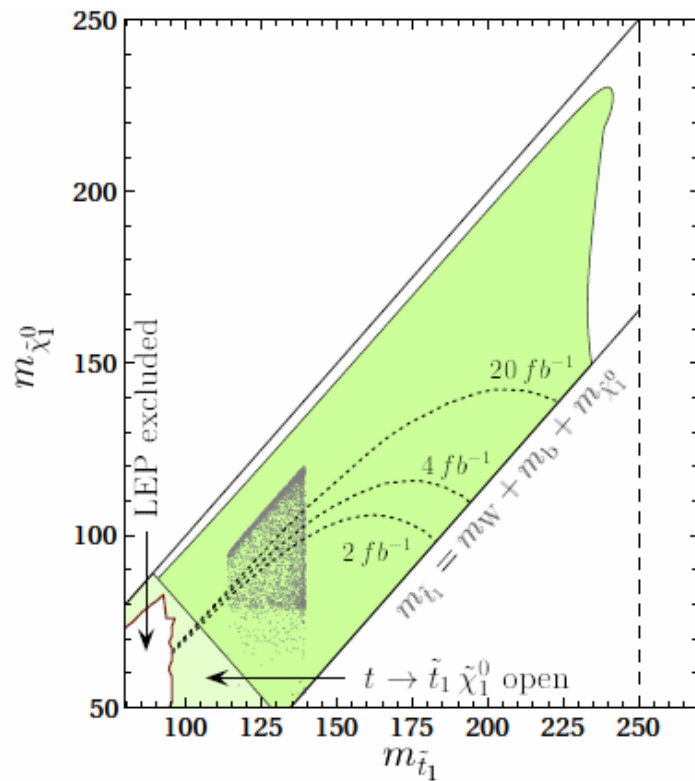
[Carena, Finch, Freitas, Milstene, Nowak, Sopczak]

[Freitas, Milstene, Schmitt, Sopczak]

[Bartl, Eberl, Kraml, Majerottom, Porod, Spoczak]

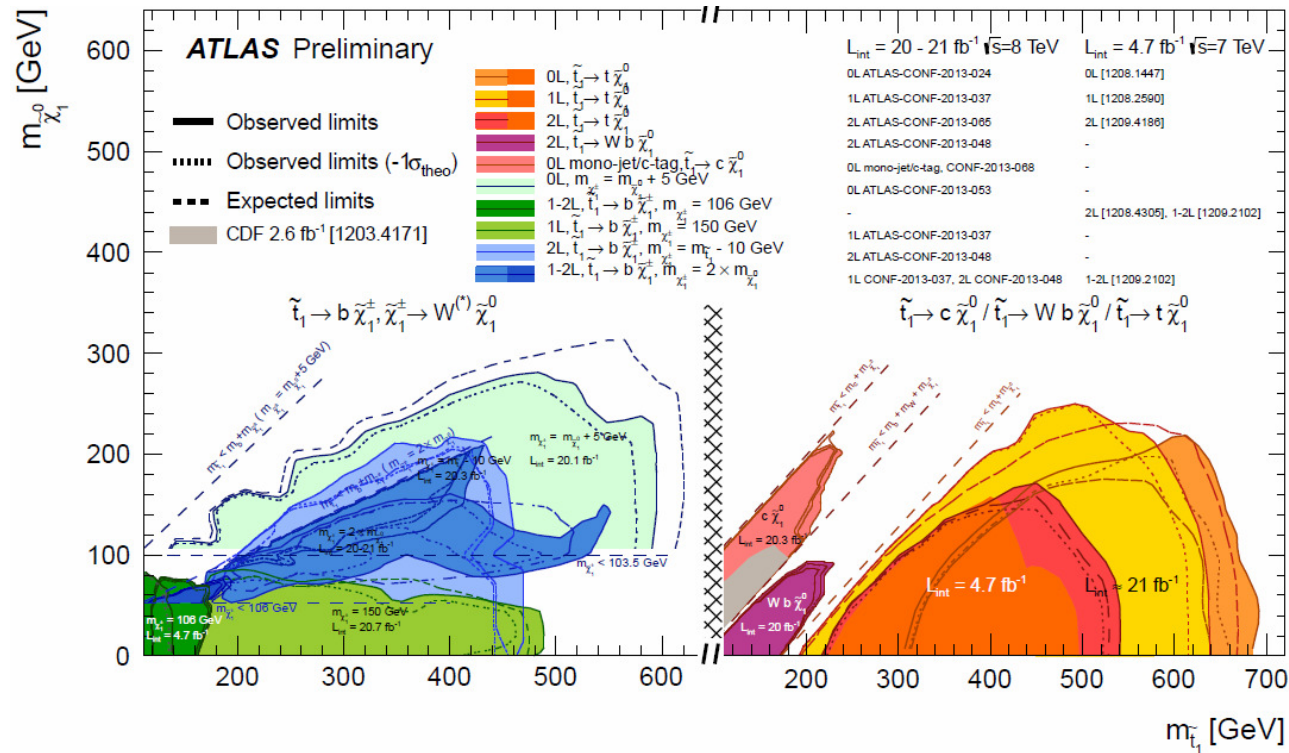
[...many others...]

New Physics at around 350 GeV: light stops



500 GeV ILC:

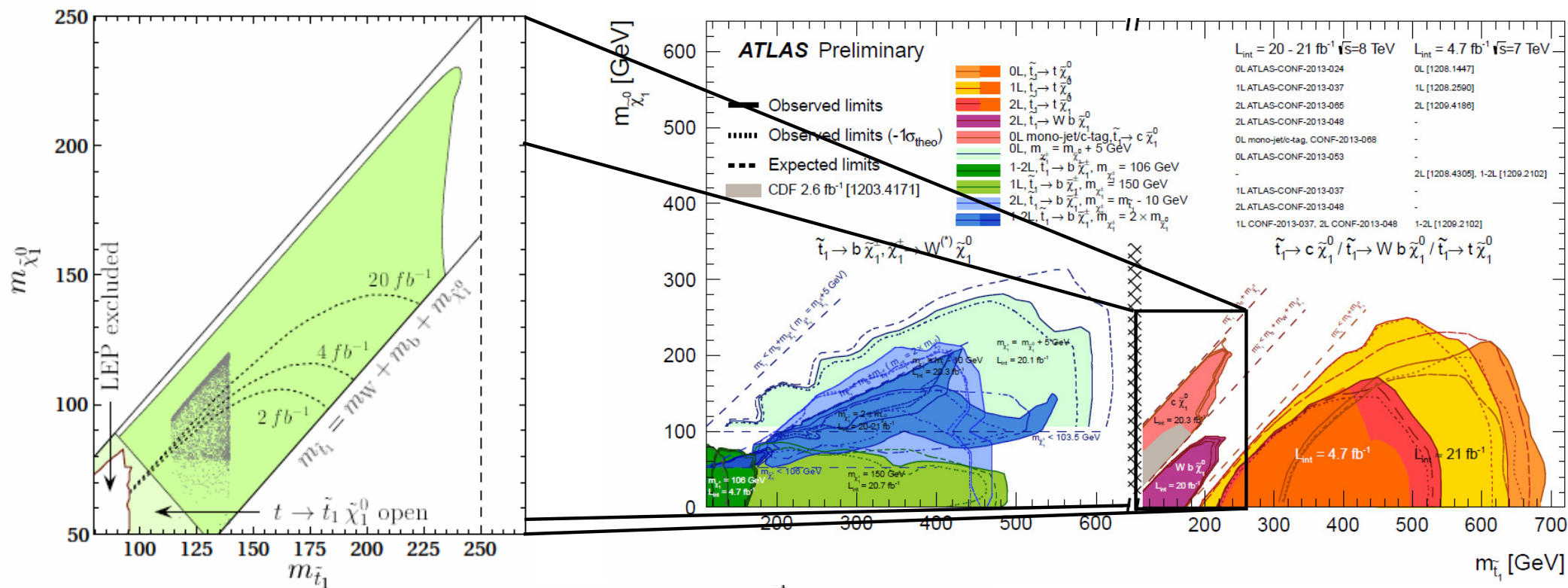
[Carena, Finch, Freitas, Milstene, Nowak, Sopczak]



7+8 TeV LHC:

[ATLAS at EPS 2013]

New Physics at around 350 GeV: light stops



500 GeV ILC:

[Carena, Finch, Freitas, Milstene,
Nowak, Sopczak]

7+8 TeV LHC:

[ATLAS at EPS 2013]

SUMMARY: Top quark physics at LC

- Future NNNLO results will most likely reduce the theoretical uncertainty of the threshold cross section to $\sim 3\%$.
- This level of accuracy is mandatory for precision top physics at threshold. It translates into an accuracy of 40-80 MeV on the top quark mass and about 2% on the top quark width.
- Realistic simulations incl. detector and beam effects, background,... do not compromise these conclusions.
- Precision on top quark electroweak couplings can be improved by an order of magnitude wrt. to LHC studies.
- Better than 20% accuracy on top quark Yukawa coupling is only possible at LC with $E \geq 500$ GeV.
- New Physics coupling to top quarks can be discovered in FCNC of the top quark decay and in single top processes.

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Realistic studies using TLEP energies and beam parameter

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What is the actual accuracy at TLEP?

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- New Physics coupling to top quarks can be discovered in FCNC of the top quark decay and in single top processes.

- Future NNNLO results will most likely reduce the theoretical uncertainty of the threshold cross section to $\sim 3\%$.

Realistic studies using TLEP energies and beam parameter

- This level of accuracy is mandatory for precision top physics at threshold. It translates into an accuracy of 40-80 MeV on the top quark mass and about 2% on the top quark width.

Comparison of ILC and TLEP projections

- Realistic simulations incl. detector and beam effects, background,... do not compromise these conclusions.

What is the actual accuracy at TLEP?

- Precision improved by an order of magnitude wrt. to LHC studies.

TLEP allows studies of $t\bar{t}$ at 350 GeV and single top at 250 GeV!

- Better possible only
- New Physics coupling to top quarks can be discovered in FCNC of the top quark decay and in single top processes.

EXTRAS

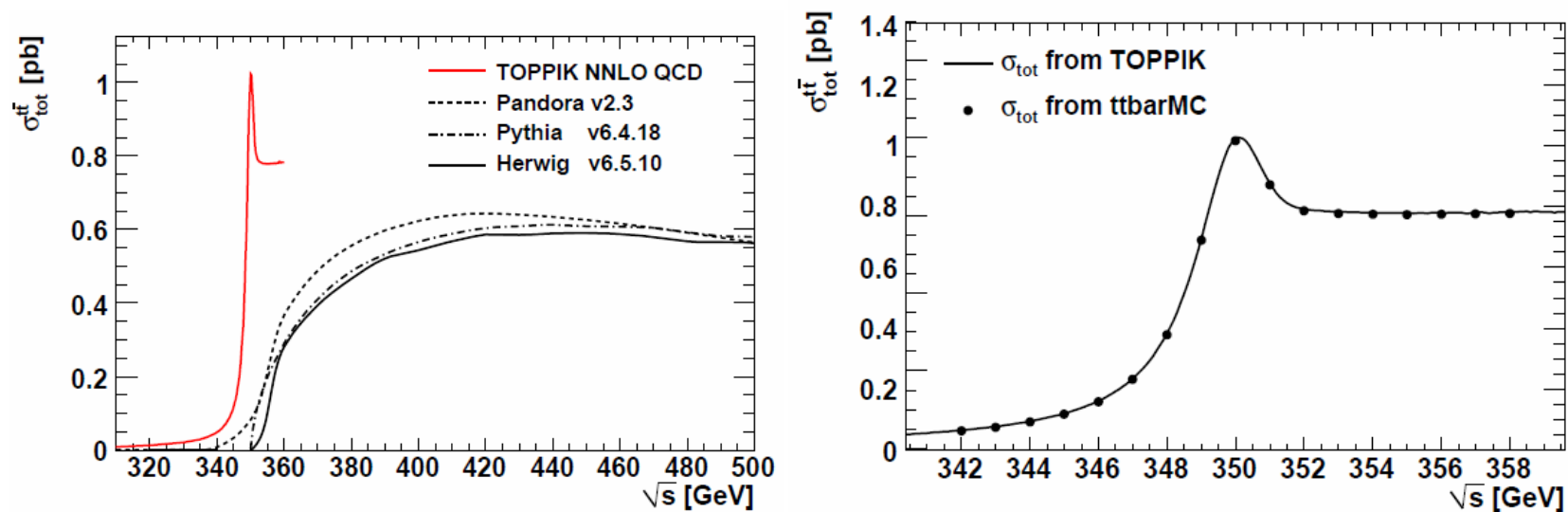
Simulation tools

MC tools for threshold production beyond LO:

TOPPIK NNLO QCD [Hoang, Teubner]

ttbarMC NNLO QCD (based on TOPPIK) [Gounaris]

CALVIN: $e^+ e^- \rightarrow \text{stops}$ (NLO SUSY QCD + coulomb corrections)
[Eberl, Bartl, Majerotto]



Extras

coupling	LHC, 300 fb ⁻¹	e^+e^- [19]
$\Delta\tilde{F}_{1V}^\gamma$	+0.043 -0.041	+0.047 , 200 fb ⁻¹ -0.047 , 200 fb ⁻¹
$\Delta\tilde{F}_{1A}^\gamma$	+0.051 -0.048	+0.011 , 100 fb ⁻¹ -0.011 , 100 fb ⁻¹
$\Delta\tilde{F}_{2V}^\gamma$	+0.038 -0.035	+0.038 , 200 fb ⁻¹ -0.038 , 200 fb ⁻¹
$\Delta\tilde{F}_{2A}^\gamma$	+0.16 -0.17	+0.014 , 100 fb ⁻¹ -0.014 , 100 fb ⁻¹
$\Delta\tilde{F}_{1V}^Z$	+0.43 -0.83	+0.012 , 200 fb ⁻¹ -0.012 , 200 fb ⁻¹
$\Delta\tilde{F}_{1A}^Z$	+0.14 -0.14	+0.013 , 100 fb ⁻¹ -0.013 , 100 fb ⁻¹
$\Delta\tilde{F}_{2V}^Z$	+0.38 -0.50	+0.009 , 200 fb ⁻¹ -0.009 , 200 fb ⁻¹
$\Delta\tilde{F}_{2A}^Z$	+0.50 -0.51	+0.052 , 100 fb ⁻¹ -0.052 , 100 fb ⁻¹

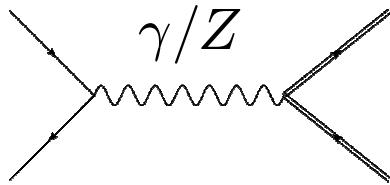
[American LC Working group]
[Baur,Juste,Orr,Rainwater]

Extras

$$e^+e^- \rightarrow t\bar{t}$$

[Devetak, Nomerotski, Peskin] (2011)

Study A_{FB} with polarized beams to determine t-tb-Z couplings

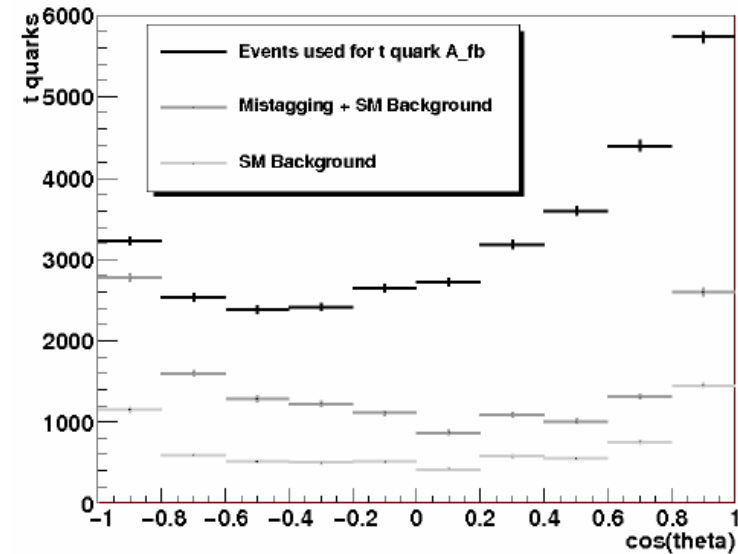


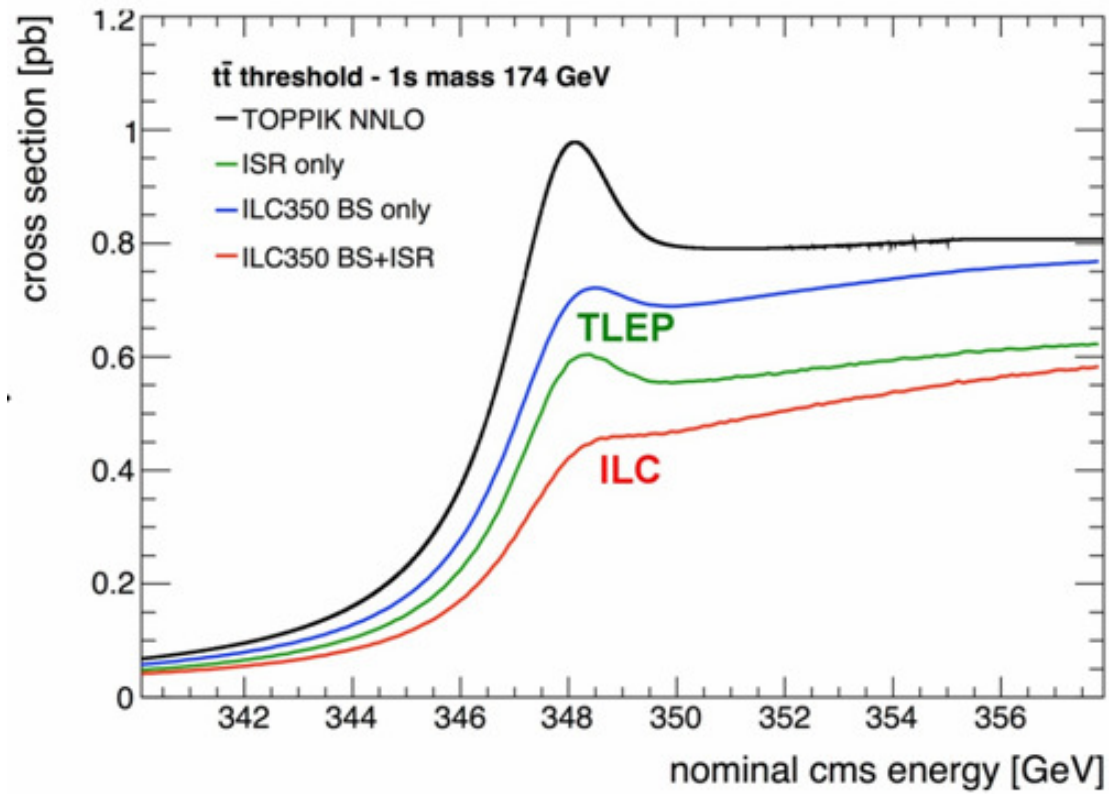
→ LO asymmetry $\sim 35\%$

- 1% precision on A_{FB} is achievable in fully hadr. channel (requires two b-tags + b charge tagging)
- polarized beams at CM energy of 500 GeV and 500 fb^{-1} :

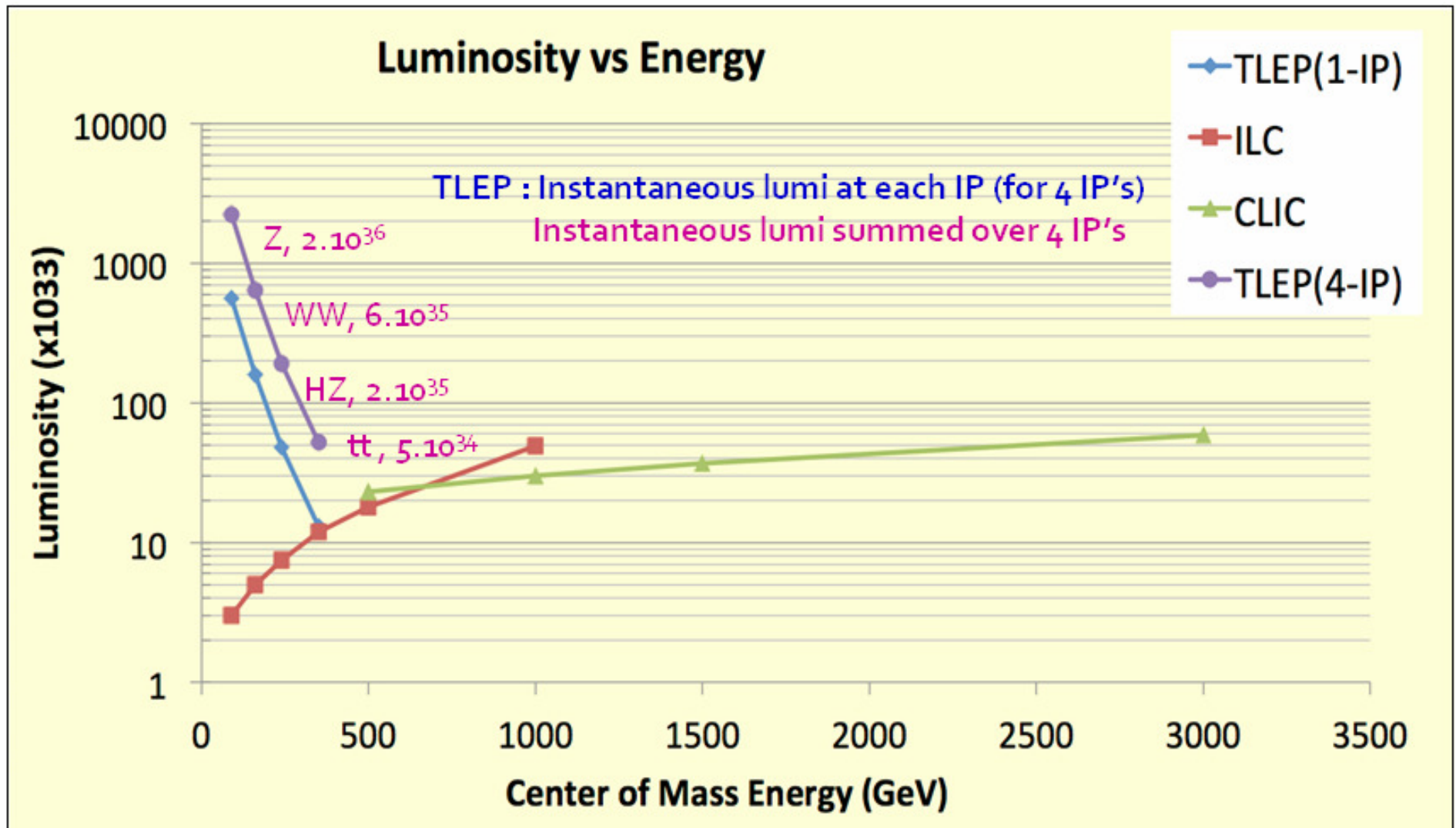
$$\delta F_L^Z \approx 6\% \quad \delta F_R^Z \approx 12\%$$

allows to place strong limits on New Physics models





[Alan Blondel, snowmass talk]



[Alan Blondel, snowmass talk]

Extras

List of Known Corrections

- QCD Source of large theory uncertainty even today(NNNLO)
 - NNLO TopWGR(2000); Hoang-Teubner, Melnikov-Yelkhovsky, Penin-Pipovarov, Beneke-Signer-Smirnov, Yakovlev, Nagano-Ota-Sumino
 - NNNLO'/NNLL' Beneke-YK-Penin-Schuller(2008), Maquard-Pichum-Seidel-Steinhauser(2006)/Hoang-Manohar-Stewart-Teubner(2001), Pineda-Signer(2006)
- EW
 - EW 1-loop Grzadkowski-Kuhn-Krawczyk-Stuart(1987), Hoang-Reisser(2006)
 - Higgs/Z-gluon 2-loop Eiras-Steinhauser(2006)
 - W-gluon 2-loop vertex YK-Seidel-Steinhauser(2008)
 - unstable top effect($t \rightarrow bW$) Hoang-Reisser-Femenia(2010)
- Susy/THDM
 - 1-loop Hollik-Schappacher(1999), Su-Wise(2001), YK-Steinhauser-Zerf(2009) /Denner-Guth-Kuhn(1992)

[Talk from Kiyo]